

Last Time:

$$\vec{\nabla} = \hat{x} \frac{\partial}{\partial x} + \hat{y} \frac{\partial}{\partial y} + \hat{z} \frac{\partial}{\partial z}$$

Gradient

$$\vec{\nabla} T$$

Divergence

$$\vec{\nabla} \cdot \vec{V}$$

Curl

$$\vec{\nabla} \times \vec{V}$$

Today: Product Rule for $\vec{\nabla}$ operations
 (Griffiths 1.2.6)

(i) $\vec{\nabla}(fg)$

(ii) $\vec{\nabla}(\vec{A} \cdot \vec{B})$

(iii) $\vec{\nabla} \cdot (\vec{f}\vec{A})$

(iv) $\vec{\nabla} \cdot (\vec{A} \times \vec{B})$

(v) $\vec{\nabla}_x (\vec{f}\vec{A})$

(vi) $\vec{\nabla}_x (\vec{A} \times \vec{B})$

6 possibilities

Griffiths front cover:

$$\nabla(fg) = f(\nabla g) + g(\nabla f)$$

$$\nabla(\mathbf{A} \cdot \mathbf{B}) = \mathbf{A} \times (\nabla \times \mathbf{B}) + \mathbf{B} \times (\nabla \times \mathbf{A}) + (\mathbf{A} \cdot \nabla)\mathbf{B} + (\mathbf{B} \cdot \nabla)\mathbf{A}$$

$$\nabla \cdot (f\mathbf{A}) = f(\nabla \cdot \mathbf{A}) + \mathbf{A} \cdot (\nabla f) \quad * \text{ Griffiths text}$$

$$\nabla \cdot (\mathbf{A} \times \mathbf{B}) = \mathbf{B} \cdot (\nabla \times \mathbf{A}) - \mathbf{A} \cdot (\nabla \times \mathbf{B}) \quad \# \text{ We'll do this one.}$$

$$\nabla \times (f\mathbf{A}) = f(\nabla \times \mathbf{A}) - \mathbf{A} \times (\nabla f)$$

$$\nabla \times (\mathbf{A} \times \mathbf{B}) = (\mathbf{B} \cdot \nabla)\mathbf{A} - (\mathbf{A} \cdot \nabla)\mathbf{B} + \mathbf{A}(\nabla \cdot \mathbf{B}) - \mathbf{B}(\nabla \cdot \mathbf{A})$$

Let's try to work on $\vec{\nabla} \cdot (\vec{A} \times \vec{B})$

$$\begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ A_x & A_y & A_z \\ B_x & B_y & B_z \end{vmatrix}$$

$$\vec{\nabla} \cdot \left[\hat{x} (A_y B_z - A_z B_y) - \hat{y} (A_x B_z - A_z B_x) + \hat{z} (A_x B_y - A_y B_x) \right]$$

$$= \frac{\partial}{\partial x} (A_y B_z - A_z B_y) - \frac{\partial}{\partial y} (A_x B_z - A_z B_x) + \frac{\partial}{\partial z} (A_x B_y - A_y B_x)$$

$$B_z \frac{\partial A_y}{\partial x} + A_y \frac{\partial B_z}{\partial x} - B_y \frac{\partial A_z}{\partial x} - A_z \frac{\partial B_y}{\partial x}$$

$$- A_x \frac{\partial B_z}{\partial y} - B_z \frac{\partial A_x}{\partial y} + A_z \frac{\partial B_x}{\partial y} + B_x \frac{\partial A_z}{\partial y}$$

$$+ A_x \frac{\partial B_y}{\partial z} + B_y \frac{\partial A_x}{\partial z} - A_y \frac{\partial B_x}{\partial z} - B_x \frac{\partial A_y}{\partial z}$$

$$B_x \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) + B_y \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right)$$

$$(\vec{\nabla} \times \vec{A})_x$$

$$(\vec{\nabla} \times \vec{A})_y$$

x -component of $\vec{\nabla} \times \vec{A}$

$$+ B_z \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) - \vec{A} \cdot (\vec{\nabla} \times \vec{B})$$

$$(\vec{\nabla} \times \vec{A})_z$$

$$\vec{B} \cdot (\vec{\nabla} \times \vec{A})$$

$$\vec{\nabla} \cdot (\vec{A} \times \vec{B}) = \vec{B} \cdot (\vec{\nabla} \times \vec{A}) - \vec{A} \cdot (\vec{\nabla} \times \vec{B})$$

1.2.7 Second Derivatives involving $\vec{\nabla}$

Possibilities are:

$$\vec{\nabla} (\vec{\nabla} \cdot \vec{v})$$

$$\vec{\nabla} \cdot (\vec{\nabla} T)$$

5 possibilities.

$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{v})$$

$$\vec{\nabla} \times (\vec{\nabla} T)$$

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{v})$$

Below, we consider the possibilities that we'll encounter in PHYS 301.

Start w/ $\vec{\nabla} \cdot (\vec{\nabla} T)$

$$\left(\hat{x} \frac{\partial}{\partial x} + \hat{y} \frac{\partial}{\partial y} + \hat{z} \frac{\partial}{\partial z} \right) \cdot \left(\frac{\partial T}{\partial x} \hat{x} + \frac{\partial T}{\partial y} \hat{y} + \frac{\partial T}{\partial z} \hat{z} \right)$$

$$= \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2}$$

This object is called the Laplacian if it is given the following notation:

$$\nabla^2 T = \vec{\nabla} \cdot (\vec{\nabla} T) = \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2}$$

Next, consider $\vec{\nabla} \times (\vec{\nabla} T)$

$$\vec{\nabla} \times (\vec{\nabla} T) = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ \frac{\partial T}{\partial x} & \frac{\partial T}{\partial y} & \frac{\partial T}{\partial z} \end{vmatrix}$$

$$= \hat{x} \left(\underbrace{\frac{\partial}{\partial y} \frac{\partial T}{\partial z} - \frac{\partial}{\partial z} \frac{\partial T}{\partial y}}_0 \right) - \hat{y} \left(\underbrace{\quad}_0 \right) + \hat{z} \left(\underbrace{\quad}_0 \right)$$

$$\frac{\partial^2 T}{\partial y \partial z} - \frac{\partial^2 T}{\partial z \partial y} = 0$$

$$\therefore \vec{\nabla} \times (\vec{\nabla} T) = 0$$

From electrostatics, we know that $\vec{E} = -\vec{\nabla} V$

$$\therefore \vec{\nabla} \times \vec{E} = \vec{\nabla} \times (-\vec{\nabla} V) = -\vec{\nabla} \times (\vec{\nabla} V) = 0$$

$$\therefore \vec{\nabla} \times \vec{E} = 0$$

One of Maxwell's
Eqns for Electrostatics

Next, consider

$$\begin{aligned}\vec{\nabla} \cdot (\vec{\nabla} \times \vec{v}) &= \vec{\nabla} \cdot \left[\hat{x} \left(\frac{\partial v_z}{\partial y} - \frac{\partial v_y}{\partial z} \right) - \hat{y} \left(\quad \right) + \hat{z} \left(\quad \right) \right] \\ &= \frac{\partial}{\partial x} \left(\frac{\partial v_z}{\partial y} - \frac{\partial v_y}{\partial z} \right) \\ &\quad - \frac{\partial}{\partial y} \left(\frac{\partial v_z}{\partial x} - \frac{\partial v_x}{\partial z} \right) \\ &\quad + \frac{\partial}{\partial z} \left(\frac{\partial v_y}{\partial x} - \frac{\partial v_x}{\partial y} \right) = 0\end{aligned}$$

Ultimately, we will see that \vec{B} can be written as $\vec{B} = \vec{\nabla} \times \vec{A}$ where \vec{A} is called a vector potential.

$$\vec{\nabla} \cdot \vec{B} = \vec{\nabla} \cdot (\vec{\nabla} \times \vec{A}) = 0$$

$$\boxed{\vec{\nabla} \cdot \vec{B} = 0}$$

Another Maxwell Eq'n.

1.3 Integral Calculus

Common integrals in $E^4 M$

Line integrals: $\Delta V = - \int_a^b \vec{E} \cdot d\vec{l}$

closed loop $\rightarrow \int \vec{B} \cdot d\vec{l} = \mu_0 I_{\text{enc}} \quad (\text{Ampere's Law})$

Surface Integrals: $\overline{\Phi} = \int \vec{B} \cdot d\vec{a} \quad (\text{flux})$

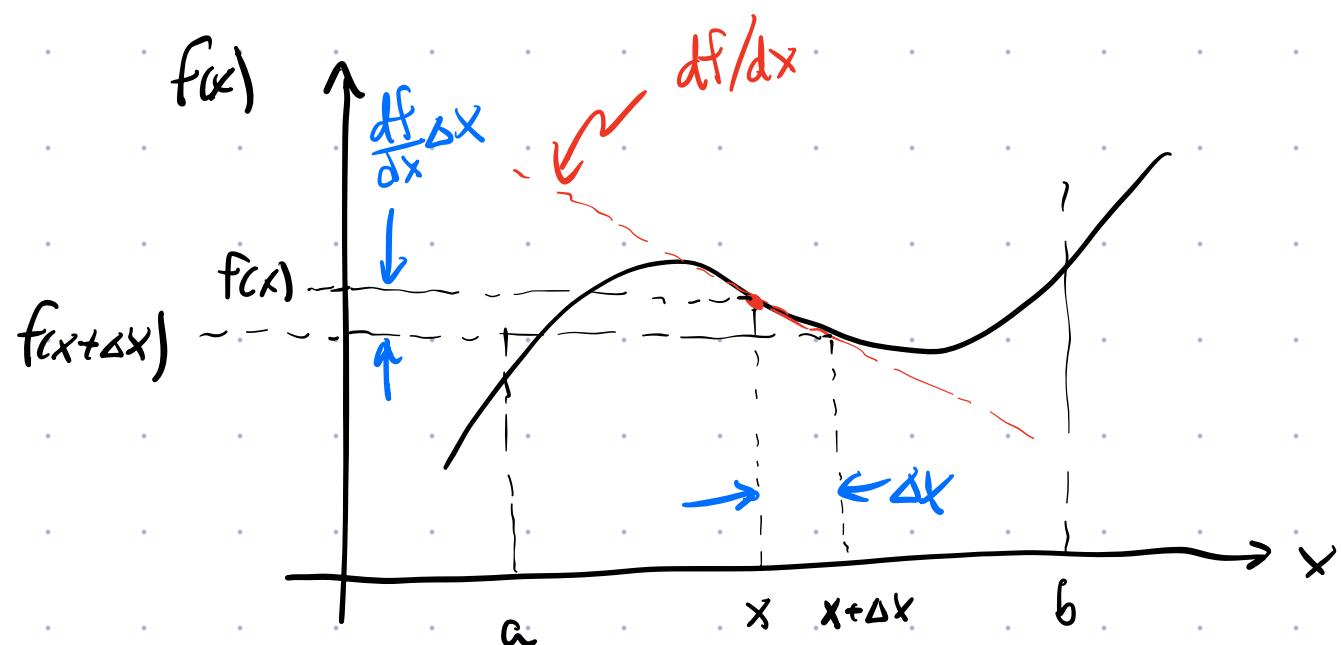
closed surface $\rightarrow \int \vec{E} \cdot d\vec{a} = \frac{q_{\text{enc}}}{\epsilon_0} \quad (\text{Gauss's Law})$

Fundamental Theorem of Calculus

$$\int_a^b \left(\frac{df}{dx} \right) dx = f(b) - f(a)$$

Re-express integral as Riemann sum

$$\int_a^b \left(\frac{df}{dx} \right) dx = \lim_{\Delta x \rightarrow 0} \sum_{x=a}^b \left(\frac{df}{dx} \right) \Delta x$$



$$\therefore \frac{df}{dx} \Delta x \approx f(x + \Delta x) - f(x)$$

$$\int_a^b \left(\frac{df}{dx} \right) dx = \lim_{\Delta x \rightarrow 0} \sum_a^b [f(x + \Delta x) - f(x)]$$

$$\begin{aligned}
 &= \left(f(a+\Delta x) - f(a) \right) + \left(f(a+2\Delta x) - f(a+\Delta x) \right) \\
 &\quad + \left(f(a+3\Delta x) - f(a+2\Delta x) \right) + \dots \\
 &\quad + \left(f(b) - f(b-\Delta x) \right)
 \end{aligned}$$

$$\boxed{\therefore \int_a^b \left(\frac{df}{dx} \right) dx = f(b) - f(a)}$$

Fundamental Theorem for Gradients

Similar to above, but applies to scalar
fns of multiple variables $\vec{T}(x, y, z)$

$$dT = \vec{\nabla} T \cdot d\vec{l} = \frac{\partial T}{\partial x} dx$$

$$+ \frac{\partial T}{\partial y} dy$$

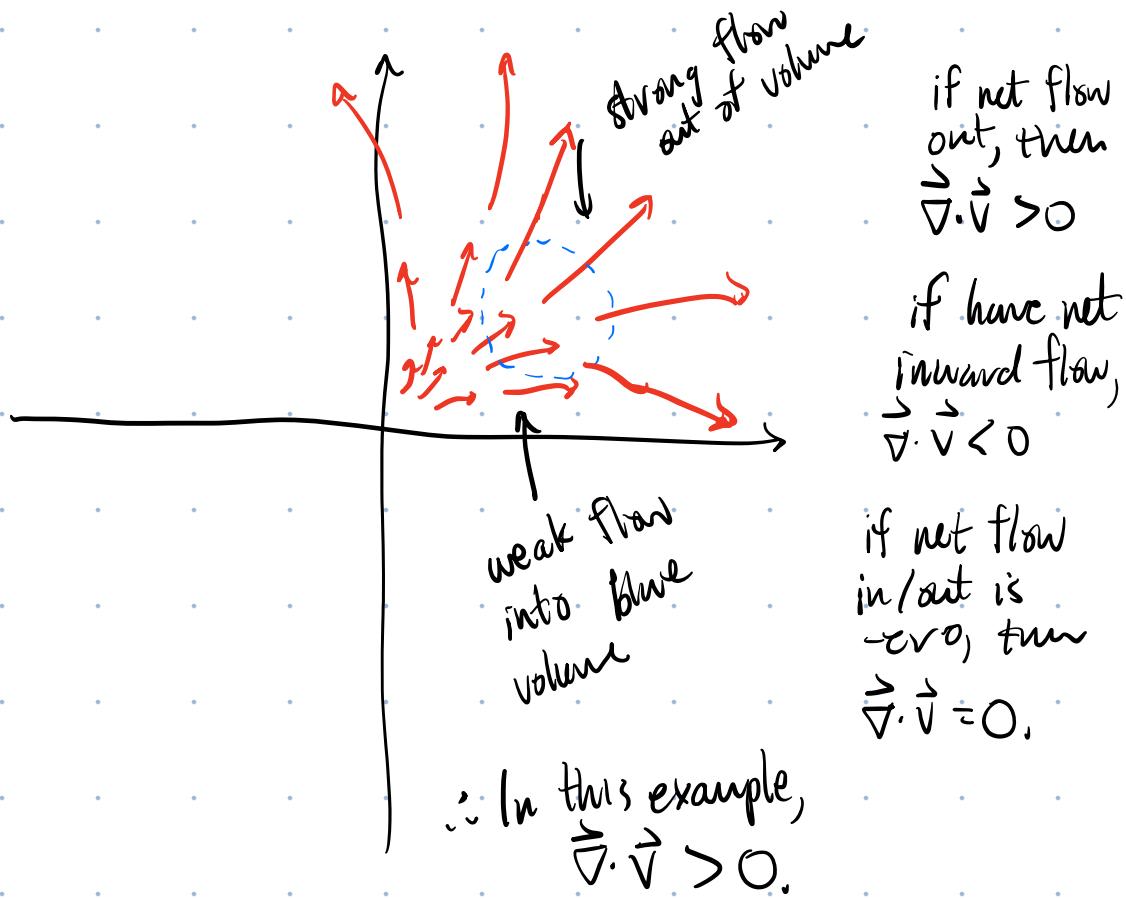
$$+ \frac{\partial T}{\partial z} dz$$

$$\int_{\vec{a}}^{\vec{b}} (\vec{\nabla} T) \cdot d\vec{l} = T(\vec{b}) - T(\vec{a})$$

pts in 3-D space

Divergence Theorem

One way to determine if a vector field \vec{V} has a non-zero gradient is to examine the "flow" of \vec{V} into/out of an infinitesimal sphere.



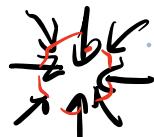
$\nabla \cdot \vec{V}$ quantifies the flow of \vec{V} in/out of a region of in space.

source or faucet of \vec{V}



$$\nabla \cdot \vec{V} > 0$$

drain or a sink of \vec{V}



$$\nabla \cdot \vec{V} < 0$$

$$\int_{\text{Volume}} \nabla \cdot \vec{V} dV \xrightarrow{dx dy dz}$$

adds up all sinks & sources of \vec{V} in a volume of space.

Some sinks & sources in the volume will cancell. If there is a net source (sink) then will have field lines exit (enter) the closed surface that surrounds the volume.

∴ We can instead simple track the flux of \vec{V} out (into) this surface.

closed surface surrounding V .

$$\oint \vec{V} \cdot d\vec{a}$$

Divergence Theorem

$$\int_V \nabla \cdot \vec{v} d\tau = \oint_S \vec{v} \cdot d\vec{a}$$

Application to electrostatics ...

Recall Gauss's Law

$$\oint_S \vec{E} \cdot d\vec{a} = \frac{q_{\text{enc}}}{\epsilon_0}$$

Apply divergence theorem:

$$\oint_S \vec{E} \cdot d\vec{a} = \int_V \nabla \cdot \vec{E} d\tau = \frac{q_{\text{enc}}}{\epsilon_0}$$

We can express $q_{\text{enc}} = \int_V \rho d\tau$

change density

$$\int_V \nabla \cdot \vec{E} d\tau = \frac{1}{\epsilon_0} \int_V P d\tau$$

$$= \int_V \left(\frac{P}{\epsilon_0} \right) d\tau$$

$\therefore \nabla \cdot \vec{E} = \frac{P}{\epsilon_0}$

another Maxwell's eqn.

Gauss's Law :

Integral form $\oint \vec{E} \cdot d\vec{a} = \frac{\text{Gauss}}{\epsilon_0}$

Differential form $\nabla \cdot \vec{E} = \frac{P}{\epsilon_0}$

Stoke's Theorem

$$\int_S (\vec{\nabla} \times \vec{v}) \cdot d\vec{a} = \oint \vec{v} \cdot d\vec{l}$$

evaluate curl of \vec{v} over surface

finds components of \vec{v} that "circulate" around boundary.

interior contribs cancel leaving only a net outer "swirl" around boundary.

