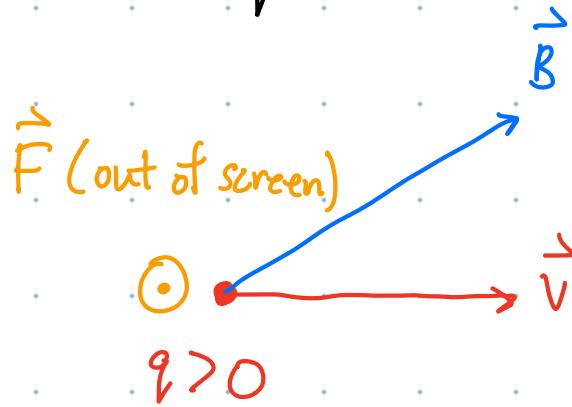


- Next PrairieLearn HW due Mar. 28
 - Complete Pre-Lab #8 before the start of your lab.
 - See course website for final exam details (including formula sheet)
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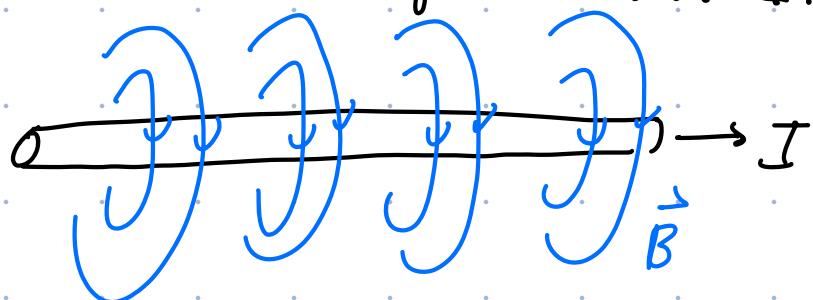
Last week : Force on a charge moving through a magnetic field.

$$\vec{F} = q \vec{v} \times \vec{B}$$



Magnetic field due to infinite straight current I .

$$B = \frac{\mu_0 I}{2\pi r}$$

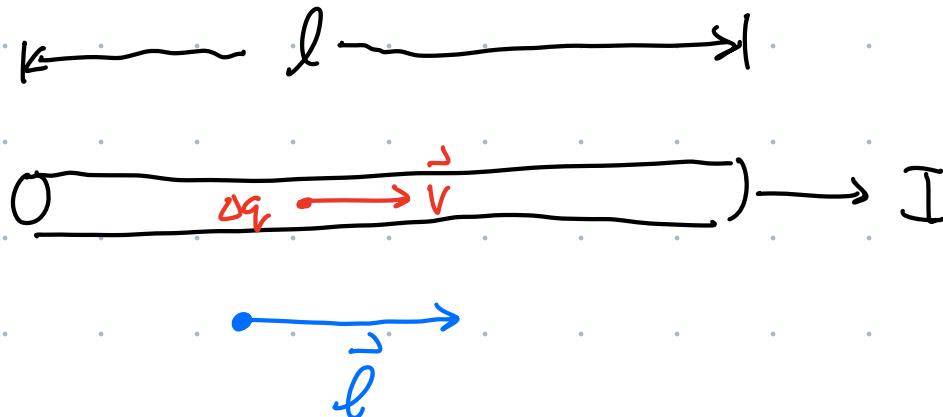


Magnetic Force on a Current.

Know for a single pt. charge in \vec{B} , the force is:

$$\vec{F} = q\vec{v} \times \vec{B} \quad ①$$

What is the force on a current (collection of moving charges) in a magnetic field?



The time Δt for a collection of charges Δq to move the length of the wire is

$$\Delta t = \frac{l}{v} \leftarrow \text{speed of charges.}$$

$$\text{Current } I = \frac{\Delta q}{\Delta t} = \frac{\Delta q}{\left(\frac{l}{v}\right)} = v \frac{\Delta q}{l}$$

$$\text{or: } \Delta q v = Il$$

define \vec{l} s.t. its dirn is given by the dirn of the current I .

$$\Delta q \vec{v} = I \vec{l}$$

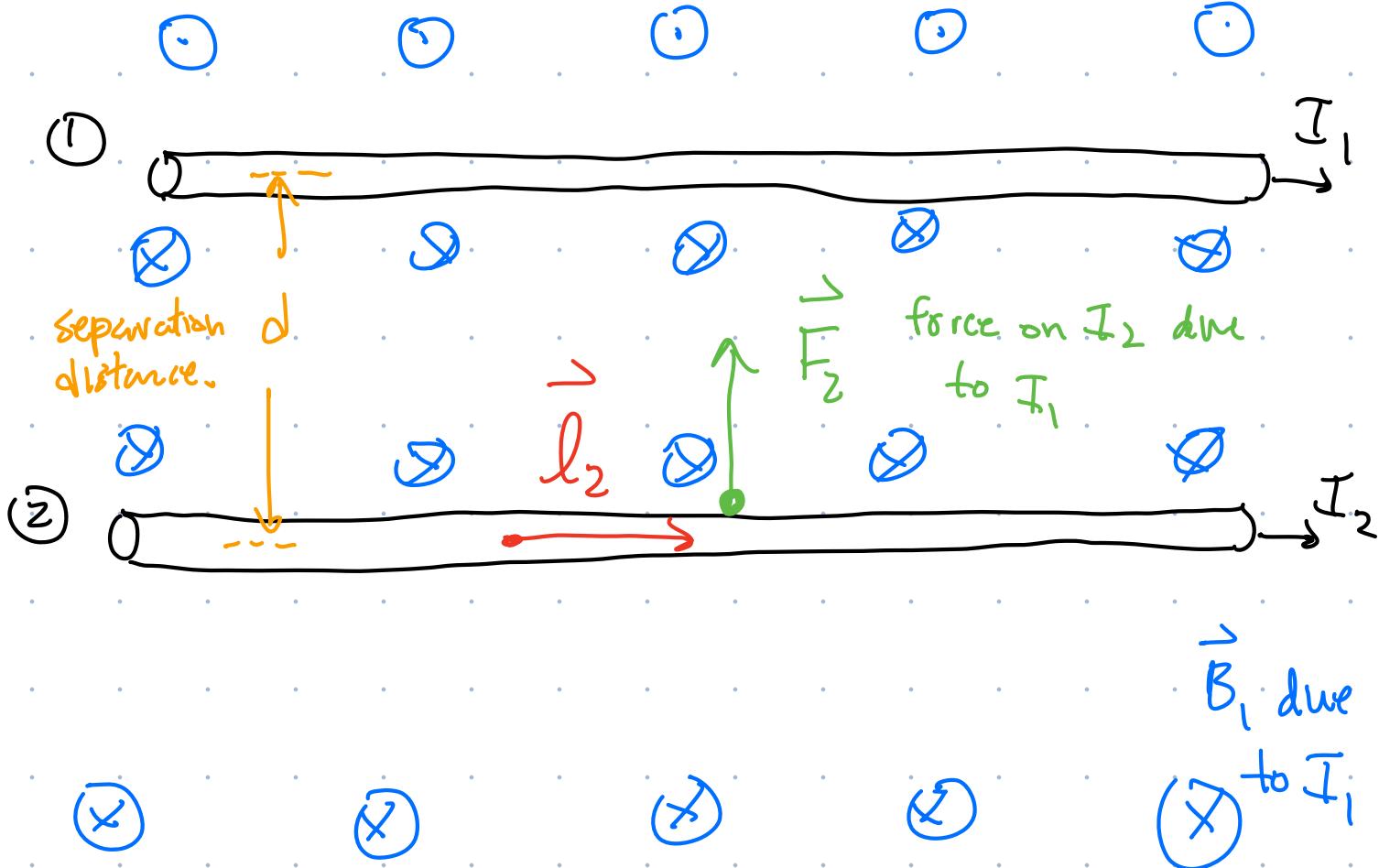
sub. into ①

$$\vec{F} = q \vec{v} \times \vec{B} \quad ①$$

$$\therefore \vec{F} = I \vec{l} \times \vec{B}$$

Force on a current in magnetic field \vec{B} .

Example: Force between Parallel Currents.



What is the force on I_2 due to magnetic field of I_1 ?

I_1 makes a magnetic field \vec{B}_1 that is into the screen at the position of I_2 .

Force on I_2 is given by: $\vec{F}_2 = I_2 \vec{l}_2 \times \vec{B}_1$

By RHR, $\vec{l}_2 \times \vec{B}_1$ is towards top of the screen.

By Newton's 3rd Law, force on I_1 due to the magnetic field of I_2 is down.

\Rightarrow Parallel currents attract.

Using the same logic, can deduce that antiparallel currents repel.

Let's work out the magnitude of $\vec{F}_2 = I_2 \vec{l}_2 \times \vec{B}_1$

$$\vec{l}_2 \perp \vec{B}_1 \Rightarrow \theta = 90^\circ$$

$$|\vec{F}_2| = I_2 |\vec{l}_2 \times \vec{B}_1| = I_2 l_2 B_1 \sin 90^\circ$$

$$|\vec{F}_2| = I_2 l B_1$$

Since I_1 is a long, straight current, know

$$B_1 = \frac{\mu_0 I_1}{2\pi d}$$

dist. from I_1 to I_2

$$\therefore |\vec{F}_2| = I_2 l \left(\frac{\mu_0 I_1}{2\pi d} \right)$$

$$\boxed{|\vec{F}_1| = |\vec{F}_2| = \frac{\mu_0 I_1 I_2 l}{2\pi d}}$$

} attractive if $I_1 \parallel I_2$
 repulsive if $I_1 \nparallel I_2$

Ampère's Law

- Gives us a way to calc. \vec{B} due to certain symmetric current distributions.
- Analogue of Gauss's Law, but for magnetic fields.

We will find that:

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{\text{enclosed}}$$

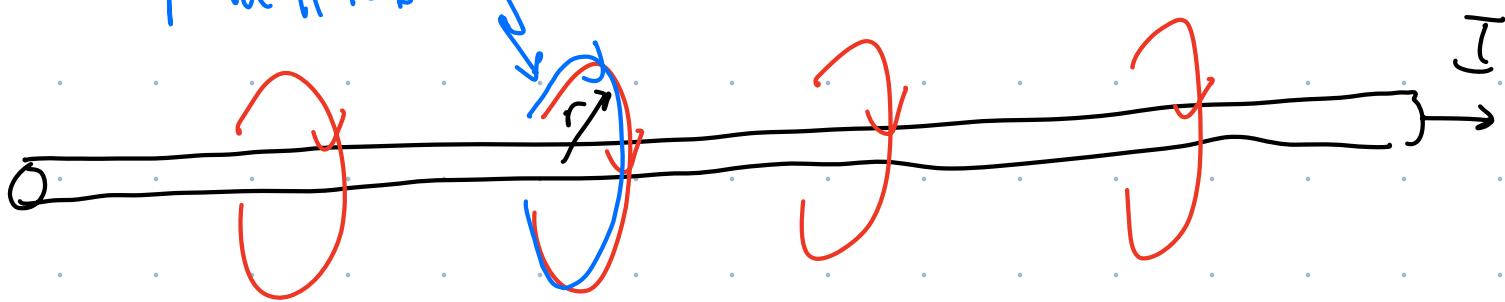
integrate
around a
closed path/loop.

Ampère's Law

Recall that \vec{B} due to a long, straight current is given by

$$B = \frac{\mu_0 I}{2\pi r}$$

Integration path \parallel to \vec{B}



Let's evaluate $\oint \vec{B} \cdot d\vec{l}$ for this long straight current. Integrate around a loop that is always parallel to \vec{B} .

$$\text{In this case, } \vec{B} \cdot d\vec{l} = B dl \cos \theta$$

$\underbrace{}_1$

$$= B dl$$

$$\oint \vec{B} \cdot d\vec{l} = \int B dl$$

For our chose path,
dist. from I to the
loop is always $r \Rightarrow B$ is
const. on our path.

$$= B \underbrace{\int dl}_{\text{circumference of path } (2\pi r)} = 2\pi r B$$

$$\therefore \oint \vec{B} \cdot d\vec{l} = 2\pi r B$$

Already know from Biot-Savart law that $B = \frac{\mu_0 I}{2\pi r}$ for a long straight current

$$\therefore \oint \vec{B} \cdot d\vec{l} = 2\pi r \left(\frac{\mu_0 I}{2\pi r} \right)$$

$$\therefore \oint \vec{B} \cdot d\vec{l} = \mu_0 I_{\text{end}}$$

We've recovered Ampère's Law.

↑ current that passes through the integration path, usually labelled as I_{end} .

What about an integration path of arbitrary shape?
Does Ampère's Law still hold?

End view of a long straight current.

Can form any arbitrary path around I using small radial sections of circular arcs.

