

To do:

- complete HW4 on PL by Jan. 31 @ 23:59

- Pre-Lab #2 is **optional**. It introduces computational methods that can be used to solve problems in physics.

- Complete it if interested / curious.

It is not worth marks, but we will grade the submissions received.

- Quiz #1 will be on Wed. Jan. 29

- See course website for details, including the formula sheet.

Last Time: Electric flux through a closed surface:

$$\Phi = \oint \vec{E} \cdot d\vec{A}$$

- ① any \vec{E} due to charges outside a closed surface contribute zero flux, Φ .
- ② $\Phi > 0$ when surface encloses $q > 0$
 $\Phi < 0$ " " " $q < 0$
- ③ Φ does not depend on position of q inside closed surface
- ④ Φ does not depend on shape of closed surface
- ⑤ Φ is proportional to the value of q enclosed by the surface,

The above observations led to Gauss's Law:

$$\therefore \vec{\Phi} = \oint \vec{E} \cdot d\vec{A} = \frac{Q_{\text{encl}}}{\epsilon_0}$$

closed surface

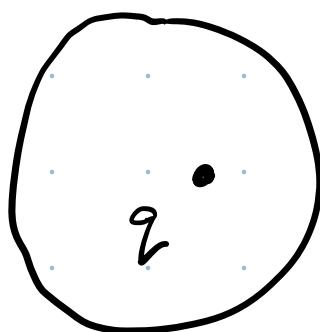
Gauss's
Law.

"encl" subscript
reminds us that
it is only the
charge inside the
closed surface that
matters.

Today: Applications of Gauss's law

E.g. What is $\vec{\Phi}$?

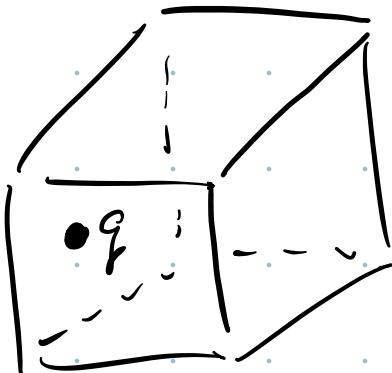
(a)



$$Q_{\text{encl}} = q$$

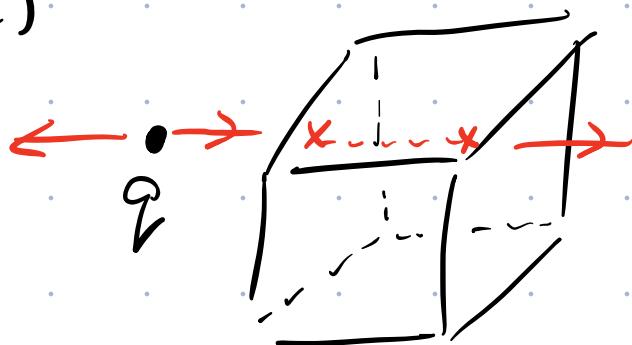
$$\therefore \vec{\Phi} = \frac{Q_{\text{encl}}}{\epsilon_0} = \frac{q}{\epsilon_0}$$

(b)



$$\Phi = \frac{q}{\epsilon_0}$$

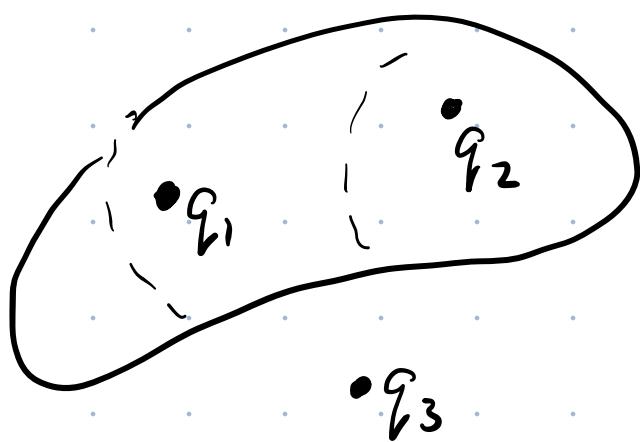
(c)



$$q_{\text{encl}} = 0$$

$$\therefore \Phi = \frac{q_{\text{encl}}}{\epsilon_0} = 0$$

(d)



$$q_{\text{encl}} = q_1 + q_2$$

(q_3 is outside surface)

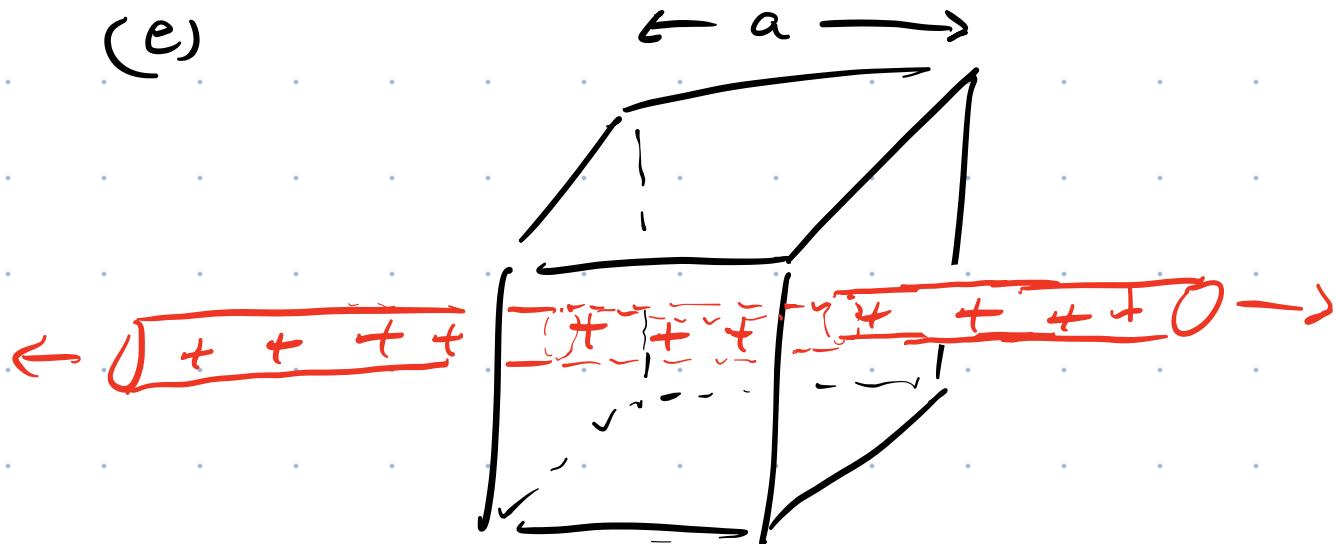
$$\therefore \Phi = \frac{q_{\text{encl}}}{\epsilon_0} = \frac{q_1 + q_2}{\epsilon_0}$$

Since we can build any distn
of charge from a collection of pt. charges,

Gauss's Law $\oint \vec{E} \cdot d\vec{A} = \frac{q_{\text{enc}}}{\epsilon_0}$ applies to all

charge dist'n (not just pt. charges).

(e)



Long, uniformly-charged rod passes through a cube of sides length a . The charge per unit length of the rod is λ . Find Φ through cube.

$$\Phi = \frac{q_{\text{enc}}}{\epsilon_0}$$

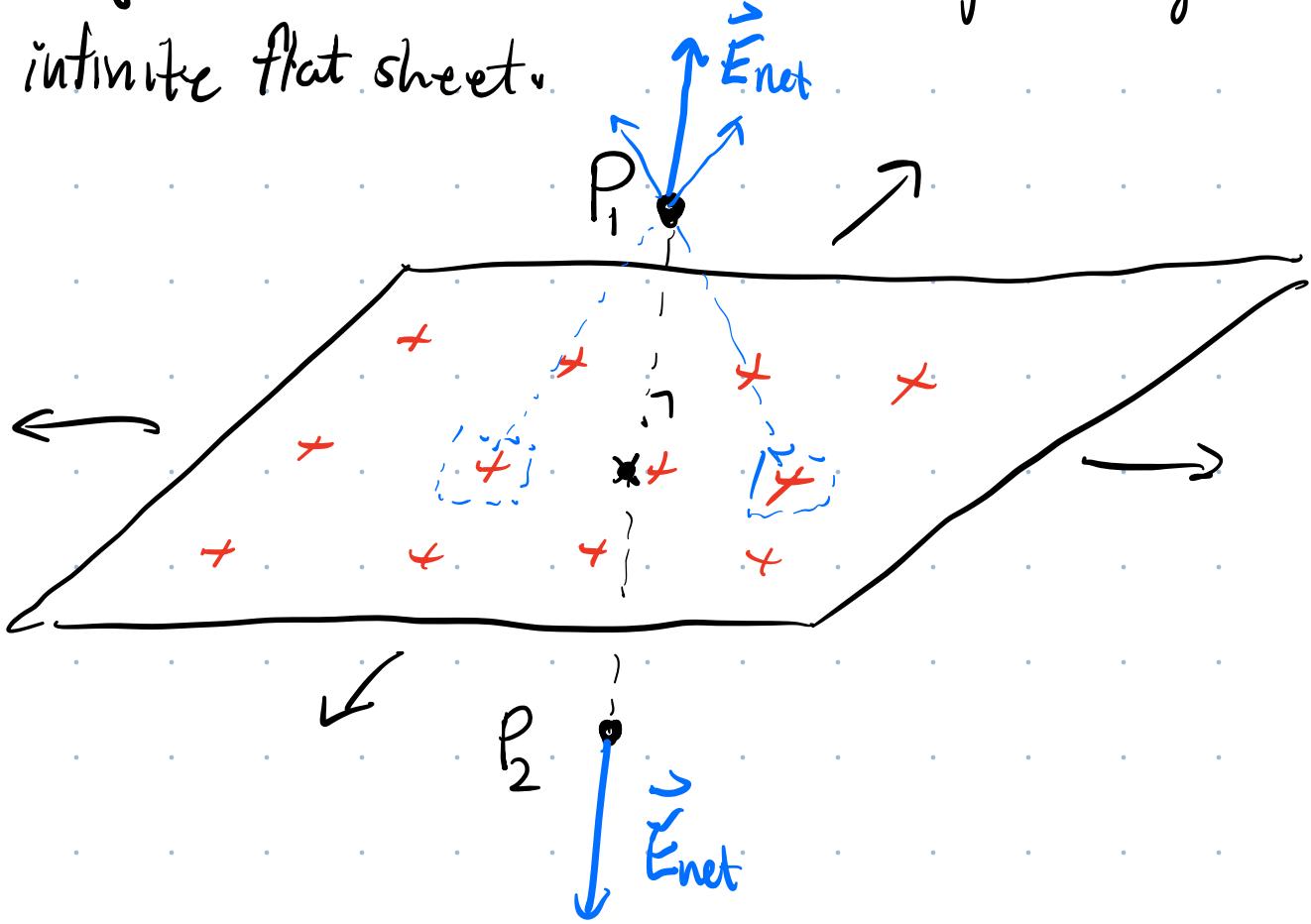
$$q_{\text{enc}} = \lambda a$$

since a length a of the rod is inside cube.

$$\Phi = \frac{\lambda a}{\epsilon_0}$$

The main utility of Gauss's is to find electric fields due to certain continuous distributions of charge.

Eg. Find \vec{E} due to a uniformly-charged infinite flat sheet.

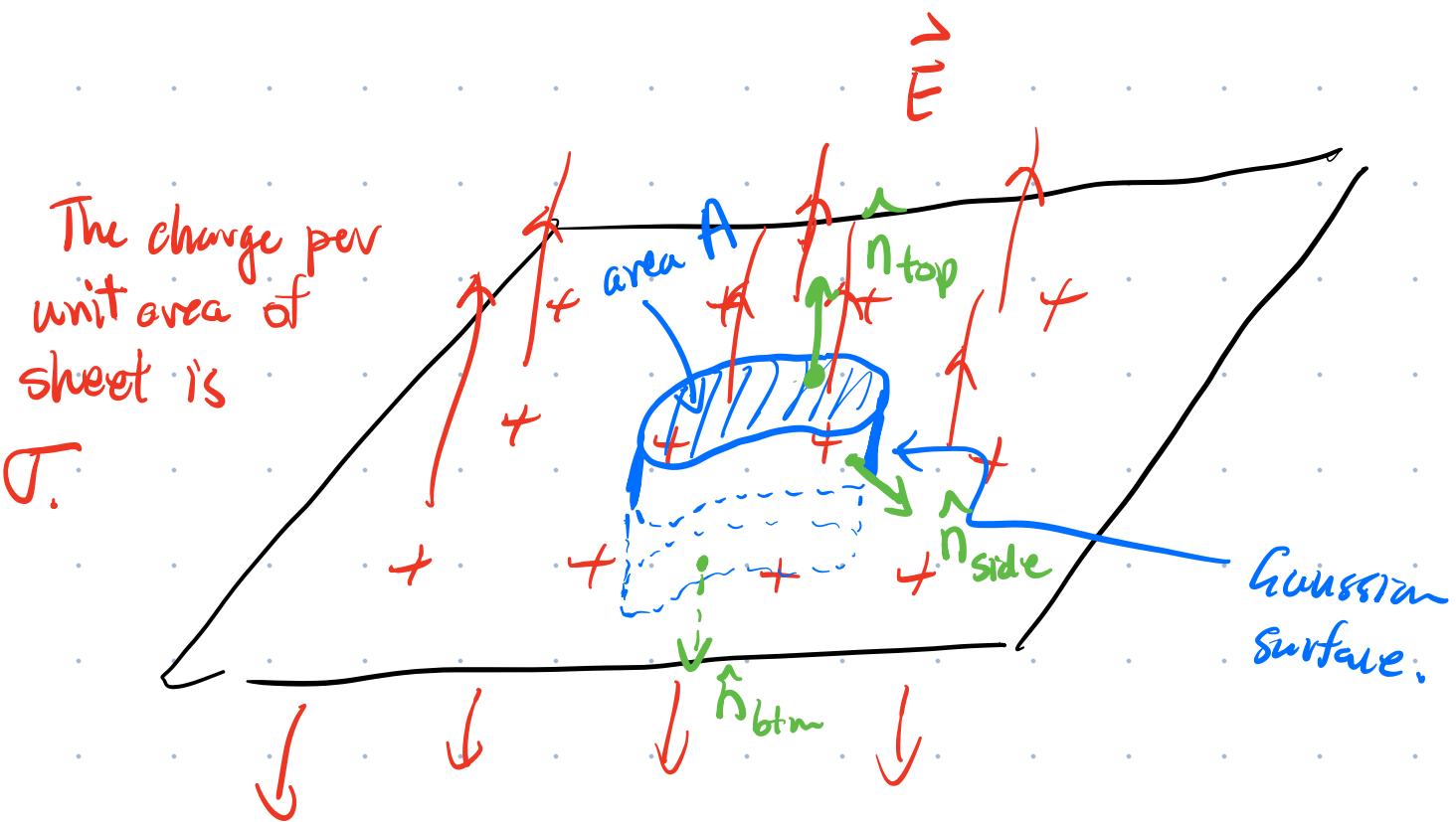


Examine symmetry of this problem before applying Gauss's law. By symmetry only the vertical contributions to \vec{E} will survive. $\therefore \vec{E}$ due to charged sheet will be \perp to the sheet.

To apply Gauss's Law, need to pick a closed surface to use to evaluate $\vec{Q} = \oint \vec{E} \cdot d\vec{A}$.

This surface is called a Gaussian surface.

To make the calculation simple, select a surface that matches the symmetry of \vec{E} & the charge dist'n.



$$\hat{n}_{top} \parallel \vec{E}$$

$$\vec{E} \cdot d\vec{A}_{top} = E dA_{top}$$

$$\hat{n}_{btm} \parallel \vec{E}$$

$$\vec{E} \cdot d\vec{A}_{btm} = E dA_{btm}$$

$$\hat{n}_{side} \perp \vec{E}$$

$$\vec{E} \cdot d\vec{A}_{side} = 0$$

$$\oint \vec{E} \cdot d\vec{A} = \int \vec{E} \cdot d\vec{A}_{top} + \int \vec{E} \cdot d\vec{A}_{btm} + \int \vec{E} \cdot d\vec{A}_{side}$$

closed surface

$$= \int E dA_{top} + \int E dA_{btm}$$

Since the top & btm parts of our Gaussian surface are parallel to the sheet, expect E to be const. on these parts of our Gaussian surface.

$$\begin{aligned}\bar{\Phi} = \oint \vec{E} \cdot d\vec{A} &= E \int_{\text{A}} dA_{top} + E \int_{\text{A}} dA_{btm} \\ &= EA + EA = 2EA\end{aligned}$$

$$\bar{\Phi} = 2EA$$

①

Next, find $\vec{\Phi}$ via $\frac{q_{\text{enc}}}{\epsilon_0}$, where q_{enc} is the charge inside our Gaussian surface.

Since the charge per unit area of the sheet is σ & the cross-sectional area of the Gaussian surface is A :

$$q_{\text{enc}} = \sigma A$$

$$\vec{\Phi} = \frac{q_{\text{enc}}}{\epsilon_0} = \frac{\sigma A}{\epsilon_0}$$

(2)

We've now found $\vec{\Phi}$ in two ways (1, 2) and they must be equal

$$2EA = \frac{\sigma A}{\epsilon_0}$$

$$\therefore E = \frac{\sigma}{2\epsilon_0}$$

Electric field due to a uniformly-charged sheet. It is indep. of distance from the sheet.

Summary:

0-D

Pt. charge $q.$

$$\bar{E} = \frac{k_e q}{r^2}$$

$$E \propto \frac{1}{r^2}$$

1-D

Line of charge $\lambda.$

$$E = \frac{\lambda}{2\pi\epsilon_0 r}$$

$$E \propto \frac{1}{r}$$

2-D

Sheet of charge σ

$$E = \frac{\sigma}{2\epsilon_0}$$

$$E \propto \frac{1}{r_0}$$

